

SHOW OFF TIME FOR GALERKIN APPROACH!

So, we have a weak formulation, i.e. $B(u, v) = F(v)$
with $u|_{\Gamma_D} = u_0 = 0$ (Γ_D is the Dirichlet boundary)

$$\text{Thus, } B(u - u_h, u - u_h) = B(u, u) + B(u_h, u_h) - 2B(u, u_h)$$

Since orthogonality of error gives $B(u - u_h, u_h) = 0$ (with u_h as a SPECIAL choice of U_h).

$$\text{Thus } \boxed{B(u, u_h) = B(u_h, u_h)} \quad \text{or } B(\underbrace{u - u_h}_{e_h}, \underbrace{u - u_h}_{e_h}) = \|e_h\|_{\Omega}^2$$

$$B(u, u) - B(u_h, u_h) \geq 0$$

$$\Rightarrow B(u, u) \geq B(u_h, u_h).$$

Now let \bar{u}_h be "0" solution obtained from "some" method such that $\bar{u}_h(x) = \sum_{i=0}^N \bar{a}_i \phi_i(x)$. Then $\tilde{e}_h = u - \bar{u}_h = u - u_h + u_h - \bar{u}_h$

$$\Rightarrow \tilde{e}_h = e_h + e_h^* \quad [e_h = u - u_h; e_h^* = u_h - \bar{u}_h]$$

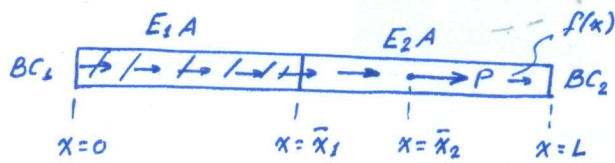
$$\text{Thus } B(\tilde{e}_h, \tilde{e}_h) = B(e_h + e_h^*, e_h + e_h^*) = B(e_h, e_h) + B(e_h^*, e_h^*) + 2B(e_h, e_h^*) = 0 \quad (\text{by orthogonality and } u_h = e_h^*)$$

$$\Rightarrow \|\tilde{e}_h\|_{\Omega}^2 = \|e_h\|_{\Omega}^2 + \|e_h^*\|_{\Omega}^2 \Rightarrow \boxed{\|\tilde{e}_h\|_{\Omega} \geq \|e_h\|_{\Omega}} \quad \text{--- (A)}$$

* Result (A) means that the Finite element solution is the CLOSEST to $u(x)$, of all the possible solutions of the form $\bar{u}_h(x)$, in ~~the~~ terms of the strain energy.

The FE/Galerkin solution is said to be the BEST APPROXIMATION of $u(x)$

Let us now look at the FE solution of a specific bar problem:

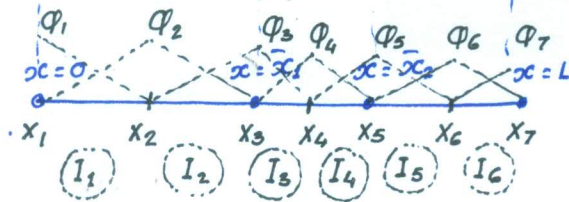


$BC_1 : u(0) = 0$ (say)

$BC_2 : E_2A \frac{du}{dx} \Big|_L = P_L$

{ A is constant }

FE mesh:



* NOTE that points of MATERIAL transition and application of point loads HAVE to be nodes.

Let $u_h(x) = \sum_{i=1}^7 \alpha_i \Phi_i(x)$ with $\alpha_1 = 0$ as $u(0) = 0$

Let $B(u_h, v_h) = F(v_h)$ with $B(u_h, v_h) = \int_0^{\bar{x}_1} E_1A \frac{du_h}{dx} \frac{dv_h}{dx} dx + \int_{\bar{x}_1}^{\bar{x}_2} E_2A \frac{du_h}{dx} \frac{dv_h}{dx} dx + \int_{\bar{x}_2}^L E_2A \frac{du_h}{dx} \frac{dv_h}{dx} dx$; $F(v_h) = \int_0^L f v_h dx + P_L v_L + P v \Big|_{\bar{x}_2}$ — (1)

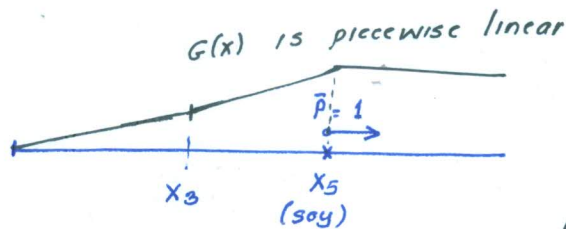
Let us also solve the problem for a different loading, given by

$f(x) = 0 ; P_L = 0 ; P = 0$ and point load $\bar{P} = 1$ at $x = x_k$ (where k is a node of the mesh). Let the exact solution of this problem be $G(x)$ and its finite element solution $G_h(x)$ with:

$G(x) \Big|_0 = G_h(x) \Big|_0 = 0 ; G_h(x) = \sum_{i=1}^7 \beta_i \Phi_i(x)$ [$\beta_1 = 0$ from $G(0) = 0$]

$\therefore G$ and G_h satisfy $B(G, v_h) = 1 \cdot v_h \Big|_{x=x_k} ; B(G_h, v_h) = 1 \cdot v_h \Big|_{x_k}$
 $\Rightarrow B(G - G_h, v_h) = 0$. — (2)

What is G ?



$G(x)$ is piecewise linear

\Rightarrow as \bar{P} is at a node,

$G_h(x) = G(x)$

[from completeness, best approximation]

$\Rightarrow B(u - u_h, v_h) = 0$ gives $B(u - u_h, G_h) = 0 = B(u - u_h, G)$.

$\Rightarrow B(G, u - u_h) = 1 \cdot (u - u_h) \Big|_{x_k} = 0 \Rightarrow \boxed{u \Big|_{x_k} = u_h \Big|_{x_k}}$ NODAL EXACTNESS!

